Simulation of QCD with a GPU cluster

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High Energy Physics

Goal: To understand the fundamental constituents of matter and their fundamental interactions



Standard Model of FUNDAMENTAL PARTICLES AND INTERACTIONS

The Standard Model is a quantum theory that summarizes our current knowledge of the physics of fundamental particles and fundamental interactions (interactions are manifested by forces and by decay rates of unstable particles)

matter constituents FERMIONS spin = 1/2, 3/2, 5/2,

tons spin =1/	2	Quark	(S spin	1 = 1/2
Mass GeV/c ²	Electric charge	Flavor	Approx. Mass GeV/c ²	Electric
(0-0.13)×10 ⁻⁹	0	up up	0.002	2/3
0.000511	-1	d down	0.005	-1/3
(0.009-0.13)×10 ⁻⁹	0	C charm	1.3	2/3
0.106	-1	S strange	0.1	-1/3
(0.04-0.14)×10 ⁻⁹	0	top	173	2/3
1.777	-1	bottom	4.2	-1/3
	tons spin = 1/ Mass GeV/c ² (0-0.13)×10 ⁻⁹ 0.000511 (0.009-0.13)×10 ⁻⁹ 0.106 (0.04-0.14)×10 ⁻⁹ 1.777	Xass GeV/c ² Electric charge (0-0.13)×10 ⁻⁹ 0 0.000511 -1 (0.009-0.13)×10 ⁻⁹ 0 0.106 -1 (0.04-0.14)×10 ⁻⁹ 0 1.777 -1	Mass GeV/c ² Electric charge Quark (0-0.13)×10 ⁻⁹ 0 up 0.000511 -1 d down (0.009-0.13)×10 ⁻⁹ 0 c c charm 0.106 -1 S strange (0.04-0.14)×10 ⁻⁹ 0 t top 1.777 -1 b bottom	Quarks spin Mass GeV/c ² Electric charge Quarks spin (0-0.13)×10 ⁻⁹ 0 Image: spin sector of the spin s

*See the neutrino paragraph below

Spin is the intrinsic angular momentum of particles. Spin is given in units of h, which is the guantum unit of angular momentum where $\hbar = h/2\pi = 6.58 \times 10^{-25}$ GeV s = 1.05×10⁻³⁴ J s.

Electric charges are given in units of the proton's charge. In SI units the electric charge of the proton is 1.60×10⁻¹⁹ coulombs.

The energy unit of particle physics is the electronvolt (eV), the energy gained by one electron in crossing a potential difference of one volt. Masses are given in GeV/c2 (remember E = mc²) where 1 GeV = 10^9 eV = 1.60×10^{-10} joule. The mass of the proton is 0.938 GeV/c² = 1.67×10⁻²⁷ kg.

Neutrinos

Neutrinos are produced in the sun, supernovae, reactors, accelerator collisions, and many other processes. Any produced neutrino can be described as one of three neutrino flavor states ν_{θ} , ν_{μ} , or ν_{τ} , labelled by the type of charged lepton associated with its production. Each is a defined quantum mixture of the three definite mass neutrinos ν_L , ν_M , and ν_H for which currently allowed mass ranges are shown in the table. Further exploration of the properties of neutrinos may yield powerful clues to puzzles about matter and antimatter and the evolution of stars and galaxy structures.

Matter and Antimatter

For every particle type there is a corresponding antiparticle type, denoted by a bar over the particle symbol (unless + or - charge is shown). Particle and antiparticle have identical mass and spin but opposite charges. Some electrically neutral bosons (e.g., Z^0 , γ , and $\eta_c = c\bar{c}$ but not $K^0 = d\bar{s}$) are their own antiparticles.

Particle Processes

These diagrams are an artist's conception. Blue-green shaded areas represent the cloud of gluons





Properties of the Interactions

tive to the strength of the electro The strengths of the interactions (forces) are shown in gnetic force for two u quarks separated by the specified distances.

Property	Gravitational Interaction	Weak Interaction (Electr	Electromagnetic Interaction	Strong Interaction
Acts on:	Mass – Energy	Flavor	Electric Charge	Color Charge
Particles experiencing:	All	Quarks, Leptons	Electrically Charged	Quarks, Gluons
Particles mediating:	Graviton (not yet observed)	W+ W- Z ⁰	γ	Gluons
Strongth at f 10 ⁻¹⁸ m	10-41	0.8	1	25
Strength at { 3×10 ⁻¹⁷ m	10-41	10-4	1	60

force carriers BOSONS spin = 0, 1, 2, ...



Strong (color) spin =1 Mass Electric Name GeV/c² charge g 0 0 gluon

Color Charge

Only quarks and gluons carry "strong charge" (also called "color charge") and can have strong interactions. Each quark carries three types of color charge. These charges have nothing to do with the colors of visible light. Just as electricallycharged particles interact by exchanging photons, in strong interactions, color-charged particles interact by exchanging gluons.

Quarks Confined in Mesons and Baryons

Quarks and gluons cannot be isolated - they are confined in color-neutral particles called hadrons. This confinement (binding) results from multiple exchanges of gluons among the color-charged constituents. As color-charged particles (quarks and gluons) move apart, the energy in the color-force field between them increases. This energy eventually is converted into additional guark-antiquark pairs. The guarks and antiquarks then combine into hadrons; these are the particles seen to emerge

Two types of hadrons have been observed in nature mesons qq and baryons qqq. Among the many types of baryons observed are the proton (uud), antiproton (uud), neutron (udd), lambda A

(uds), and omega Ω^- (sss). Quark charges add in such a way as to make the proton have charge 1 and the neutron charge 0. Among the many types of mesons are the pion π^+ (ud), kaon K⁻ (su), B⁰ (db), and n_c (cc). Their charges are +1, -1, 0, 0 respectively.





Driven by new puzzles in our understanding of the physical world, particle physicists are following paths to new wonders and

Why No Antimatter?

Matter and antimatter were created in the Big

Bang. Why do we now see only matter except

in the lab and observe in cosmic rays?

for the tiny amounts of antimatter that we make

Universe Accelerating?



The expansion of the universe appears to be accelerating. Is this due to Einstein's Cosmological Constant? If not, will experiments reveal a new force of nature or even extra (hidden) dimensions of space?

Unsolved Mysteries

startling discoveries. Experiments may even find extra dimensions of space, mini-black holes, and/or evidence of string theory.



Invisible forms of matter make up much of the mass observed in galaxies and clusters of galaxies. Does this dark matter consist of new types of particles that interact very weakly with ordinary matter?

Origin of Mass?

In the Standard Model, for fundamental particles to have masses, there must exist a particle called the Higgs boson. Will it be discovered soon? Is supersymmetry theory correct in predicting more than one type of Higgs?

Fundamental Interactions

Properties of the Interactions

The strengths of the interactions (forces) are shown relative to the strength of the electromagnetic force for two u quarks separated by the specified distances.

Property	Gravitational Interaction	Weak Electromagnetic Interaction (Electroweak)		Strong Interaction
Acts on:	Mass – Energy	Flavor	Electric Charge	Color Charge
Particles experiencing:	All	Quarks, Leptons	Electrically Charged	Quarks, Gluons
Particles mediating:	Graviton (not yet observed)	W ⁺ W ⁻ Z ⁰	γ	Gluons
Strength at $\int 10^{-18} m$	10 ⁻⁴¹	0.8	1	25
3×10 ⁻¹⁷ m	10 ⁻⁴¹	10 ⁻⁴	1	60

Fundamental Particles







Tunnel of the Large Hadron Collider



The tunnel is buried around 50 to 175 m. underground

ATLAS Detector



LHC experiments will produce roughly 15 petabytes of data annually. CERN is collaborating with institutions in 33 different countries (including Taiwan) to operate a distributed computing and data storage infrastructure known as the LHC Computing Grid



ABSTRACT

Quantum Chromodynamics (QCD) is the quantum field theory of the strong interaction, describing the interactions of the quarks and gluons making up hadrons (e.g., proton, neutron, and pion).

Quantum Chromodynamics (QCD) accounts for the nuclear energy inside an atom, and plays an important role in the evolution of the early universe.

To solve QCD is a grand challenge among all sciences. The most promising approach to solve QCD nonperturbatively is to discretize the continuum space-time into a 4 dimensional lattice (lattice QCD), and to compute physical observables with Monte Carlo simulation.

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ABSTRACT (cont)

For lattice QCD with exact chiral symmetry,
 it often requires supercomputers (e.g., 10 racks of IBM
 BlueGene supercomputer) to perform the simulations.



10 racks of IBM BG/L at KEK, Japan (57 TFLOPS peak)

ABSTRACT (cont)

TWQCD Collaboration is the first lattice QCD group around the world to use a GPU cluster (128 TFLOPS) to perform large-scale unquenched simulations of lattice QCD with exact chiral symmetry.





- Introduction
- Optimal Domain-Wall Fermion
- HMC Simulation with ODWF
- GPU Strategies
- Research Highlights
- Conclusion and Outlook

Quantum Chromodynamics (QCD)

QCD is the quantum field theory for the strong interaction, describing the interactions of the quarks and gluons making up hadrons (e.g. proton, neutron, and pion). It accounts for the nuclear energy inside the nucleus of an atom, and plays an important role in the evolution of the early universe.

Salient features:

- Gauge group $SU(3) \Rightarrow$ gluons have self-interactions.
- Asymptotic freedom: $g(r) \rightarrow 0$ as $r \rightarrow 0$.
- IR slavory: $g(r) \simeq 1$ at $r \simeq 10^{-15}$ m \Rightarrow quark/color confinement
- No exact analytic solutions

Quarks

Quarks are spin $\frac{1}{2}$ fermions carrying color, and there are 6 species (flavors) of quarks.

uctuctdsbdsbds

Hadrons are color singlets composed of quarks

P = uud + antisym. in color, Proton

N = udd + antisym. in color, Neutron

 $\pi^+ = \overline{du} + \overline{du} + \overline{du},$ Pion

The nuclear force between nucleons emerges as residual interactions of QCD

The Challenge of **QCD**

At the hadronic scale, $g(r) \simeq 1$, perturbation theory is incapable to extract any quantities from **QCD**, nor to tackle the most interesting physics, namely, the **spontaneously chiral symmetry breaking** and the **color confinement**

To extract any physical quantities from the first principles of **QCD**, one has to solve **QCD** nonperturbatively.

A viable nonperturbative formulation of **QCD** was first proposed by **K. G. Wilson** in 1974.

But, the problem of lattice fermion, and how to **formulate exact chiral symmetry on the lattice had not been resolved until 1992-98** [Kaplan, Neuberger, Narayanan,...], i.e., Lattice QCD with exact chiral symmetry.

Basic notions of Lattice QCD

1. Perform Wick rotation: $t \rightarrow -ix_4$, then $\exp(iS) \rightarrow \exp(-S_E)$, and the expectation value of any observable *O*

$$\langle O \rangle = \frac{1}{Z} \int [dA] [d\psi] [d\overline{\psi}] O (A, \psi, \overline{\psi}) e^{-S_E}$$
$$Z = \int [dA] [d\psi] [d\overline{\psi}] e^{-S_E}$$



2. Discretize the space-time as a 4-d lattice $L^4 = (Na)^4$ with lattice spacing a. Then the path integral in QFT becomes a well-defined multiple integral which can be evaluated via Monte Carlo

$$\langle O \rangle = \frac{1}{Z} \int \prod_{i} dA_{i} \prod_{j} d\psi_{j} \prod_{k} d\overline{\psi}_{k} O(A, \psi, \overline{\psi}) e^{-S_{E}}$$

Gluon fields on the Lattice

The SU(3) color gluon field $A_{\mu}(x)$ are defined on each link connecting x and $x + a\hat{\mu}$, through the link variable

$$U_{\mu}(x) = \exp\left[iagA_{\mu}\left(x + \frac{a}{2}\hat{\mu}\right)\right]$$

Then the gluon action on the lattice can be written as

$$S_{g}\left[U\right] = \frac{6}{g^{2}} \sum_{\text{plaquette}} \left[1 - \frac{1}{3} \operatorname{Re} tr\left(U_{p}\right)\right] \underline{a \to 0} \int d^{4}x \frac{1}{2} tr\left[F_{\mu\nu}\left(x\right)F_{\mu\nu}\left(x\right)\right]$$

where $U_p = U_\mu(x)U_\nu(x+a\hat{\mu})U^\dagger_\mu(x+a\hat{\nu})U^\dagger_\nu(x)$



Lattice QCD

The QCD action $S = S_G(U) + \overline{\psi}D(U)\psi$ $S_{G}(U)$ is the action of the gluon fields where $\overline{\psi}D(U)\psi \equiv \overline{\psi}_{a\alpha x}^{f}D_{f}(U)_{a\alpha x,b\beta y}\psi_{b\beta y}^{f}$ f = u, d, s, c, b, tflavor index a, b = 1, 2, 3color index $\alpha,\beta=1,2,3,4$ Dirac index $x, y = 1, \dots, N_{\text{sites}} = N_x N_y N_z N_t$ site index For example, on the $16^3 \times 32$ lattice, *D* is a complex matrix of size 1,572,864 × 1,572,864 $\left\langle \mathcal{O}(\bar{\psi},\psi,U)\right\rangle = \frac{\int dU d\bar{\psi} d\psi \mathcal{O}(\bar{\psi},\psi,U) e^{-S}}{\int dU d\bar{\psi} d\psi e^{-S}} = \frac{\int dU \Theta(D^{-1},U) \det(D) e^{-S_G}}{\int dU \det(D) e^{-S_G}}$

The Challenge of Lattice **QCD**

- To have lattice volume large enough such that m_πL ≫ 1.
 So far, the lightest u/d quark cannot be put on the lattice.
 Rely on ChPT to extrapolate lattice results to physical ones.
- To have lattice spacing small enough such that $m_a a \ll 1$.
- To meet the above two conditions:
 The lattice size should be at least 100³ × 200.
 The computing power should be at least Petaflops × year

The Hardware



One unit (1U) of Tesla \$1070



4 GPUs, each with 4 GB



# of Tesla GPUs	4	
# of Streaming Processor Cores	960 (240 per processor)	
Frequency of processor cores	1.296 to 1.44 GHz	
Single Precision floating point performance (peak)	3.73 to 4.14 TFlops	
Double Precision floating point performance (peak)	311 to 345 GFlops	
Floating Point Precision	IEEE 754 single & double	
Total Dedicated Memory	16GB	
Memory Interface	512-bit	
Memory Bandwidth	408GB/sec	
Max Power Consumption	800 W	
System Interface	PCIe x16 or x8	
Software Development Tools	<u>C-based CUDA Toolkit</u>	

Two GTX285 on one motherboard



The Hardware (cont)

- 16 units of Nvidia Tesla S1070 (total 64 GPUs, 64 x 4 GB) connected to 16 servers (total 32 Intel QC Xeon, 16 x 32 GB)
- 64 Nvidia GTX285 (total 64 GPUs, 64 x 2/1 GB), connected to 32 servers (total 32 Intel i7, 32 x 12 GB)
- Hard disk storage > 300 TB, Lustre cluster file system
- Peak performance is **128 TFLOPS**
- Developed efficient CUDA codes for unquenched lattice QCD.
 140/100 Gflops per each GTX285/T10
- Attaining 15 TFLOPS (sustained) with a price \$220,000.

Optimal Domain-Wall Fermion

$$[\text{TWC, Phys. Rev. Lett. 90 (2003) 071601 }]$$

$$A_{\text{odwf}} = \sum_{s,s'=1}^{N_s} \sum_{x,x'} \overline{\psi}_{x,s} \Big[(I + \omega_s D_w)_{x,x'} \,\delta_{s,s'} - (I - \omega_s D_w)_{x,x'} \left(P_- \delta_{s',s+1} + P_+ \delta_{s',s-1} \right) \Big] \psi_{x',s'}$$

$$\equiv \overline{\Psi} D_{\text{odwf}} \Psi \qquad D_w = \sum_{\mu=1}^4 \gamma_\mu t_\mu + W - m_0, \quad m_0 \in (0,2)$$

$$t_\mu(x,x') = \frac{1}{2} \Big[U_\mu(x) \delta_{x',x+\mu} - U_\mu^{\dagger}(x') \delta_{x',x-\mu} \Big]$$

$$W(x,x') = \sum_{\mu=1}^4 \frac{1}{2} \Big[2\delta_{x,x'} - U_\mu(x) \delta_{x',x+\mu} - U_\mu^{\dagger}(x') \delta_{x',x-\mu} \Big]$$

with boundary conditions

$$P_{+}\psi(x,0) = -\frac{m_{q}}{2m_{0}}P_{+}\psi(x,N_{s})$$
$$P_{-}\psi(x,N_{s}+1) = -\frac{m_{q}}{2m_{0}}P_{-}\psi(x,1)$$

Optimal Domain-Wall Fermion (cont.)

The weights $\{\omega_s\}$ are fixed such that the effective 4D Dirac operator possesses the optimal chiral symmetry,

$$\omega_{s} = \frac{1}{\lambda_{\min}} \sqrt{1 - \kappa'^{2} s n^{2} \left(v_{s}; \kappa' \right)}, \quad s = 1, \cdots, N_{s}$$

where $sn(v_s;\kappa')$ is the Jacobian elliptic function with argument v_s and modulus $\kappa' = \sqrt{1 - \lambda_{\min}^2 / \lambda_{\max}^2}$, λ_{\min}^2 and λ_{\max}^2 are lower and upper bounds of the eigenvalues of H_w^2

The action for Pauli-Villars fields is similar to A_{odwf}

$$A_{PV} = \sum_{s,s'=1}^{N_s} \sum_{x,x'} \bar{\phi}_{x,s} \Big[\big(I + \omega_s D_w \big)_{x,x'} \,\delta_{s,s'} - \big(I - \omega_s D_w \big)_{x,x'} \big(P_- \delta_{s',s+1} + P_+ \delta_{s',s-1} \big) \Big] \phi_{x',s'} \Big]$$

but with boundary conditions: $P_+\phi(x,0) = -P_+\phi(x,N_s)$, $m_{PV} = 2m_0$ $P_-\phi(x,N_s+1) = -P_-\phi(x,1)$

Optimal Domain-Wall Fermion (cont.)

 $\int [d\bar{\psi}] [d\psi] [d\bar{\phi}] [d\phi] \exp(-A_{\text{odwf}} - A_{\text{PV}}) = \det D(m_q)$

The effective 4D Dirac operator

D



(1847 - 1878)

$$\begin{pmatrix} m_q \end{pmatrix} = m_q + \begin{pmatrix} m_0 - m_q/2 \end{pmatrix} \begin{bmatrix} 1 + \gamma_5 S_{opt} (H_w) \end{bmatrix}$$

$$S_{opt} (H_w) = \begin{cases} H_w R_Z^{(n,n)} (H_w^2), & N_s = 2n + 1 \\ H_w R_Z^{(n-1,n)} (H_w^2), & N_s = 2n \end{cases}$$

$$I$$

$$Zolotarev optimal rational approx. for \frac{1}{\sqrt{H_w^2}}$$

$$(1877)$$

The salient feature of optimal rational approximation

Has (n+m+2) alternate change of sign in $[x_{\min}, x_{\max}]$, and attains its max. and min. (all with equal magnitude)

 $1 - \sqrt{x} R_Z^{(n,m)}(x)$

In the figure, n = m = 6, it has 14 alternate change of sign in [1,1000]



Even-Odd Preconditioning of ODWF Matrix

$$\begin{split} \left[D_{\text{odwf}} \right]_{x,s;x',s'} &= \left(I + \omega_s D_w \right)_{x,x'} \delta_{s,s'} - \left(I - \omega_s D_w \right)_{x,x'} \left(P_- \delta_{s',s+1} + P_+ \delta_{s',s-1} \right) \\ &= \left(\begin{matrix} X & D_w^{eo} Y \\ D_w^{oe} Y & X \end{matrix} \right) \\ Y_{s,s'} &= \omega_s (I + L)_{s,s'} \\ X_{s,s'} &= (4 - m_0) \omega_s (I + L)_{s,s'} + (I - L)_{s,s'} \\ L_{s,s'} &= P_- \delta_{s',s+1} + P_+ \delta_{s',s-1} \end{split}$$

with boundary conditions: $P_- L_{N_s,s'} = -\frac{m_q}{2m_0} P_- \delta_{s',1}$

$$P_{+}L_{1,s'} = -\frac{m_q}{2m_0}P_{+}\delta_{s',N_s}$$

Even-Odd Preconditioning of ODWF Matrix (cont)

 $\det D_{\text{odwf}} \Longrightarrow \det(I - D_w^{oe}YX^{-1}D_w^{eo}YX^{-1}) = \det C$

$$C \equiv I - D_w^{oe} Y X^{-1} D_w^{eo} Y X^{-1}$$

For 2-flavor QCD, the pseudofermion action is

$$A_{PF} = \phi^{\dagger} C_{PV}^{\dagger} (CC^{\dagger})^{-1} C_{PV} \phi$$

 $C_{PV} \equiv C(m_a = 2m_0)$

Hybrid Monte Carlo (HMC) for 2 flavor QCD

- 1. Initial gauge configuration $\{U_l\}$
- 2. Generate $\{P_l^a\}$ with probability distribution $\propto \exp[-(P_l^a)^2/2]$
- 3. Generate ξ with probability distribution $\propto \exp(-\xi^{\dagger}\xi)$ Recall: $\exp[-\phi^{\dagger}C_{PV}^{\dagger}(CC^{\dagger})^{-1}C_{PV}\phi] = \exp[-\xi^{\dagger}\xi]$
- 4. Fixing the pseudofermion field $\phi = C_{PV}^{-1}C\xi \equiv D\xi$
- 5. Molecular dynamics (Omelyan integrator with multiple-time scale)

 \$\begin{aligned}
 \$\exists(\alpha) = \begin{bmatrix} DD^{\dagge}(U(\tau)) \begin{aligned}
 \$-1\$ when most expensive part of HMC

 \$\begin{aligned}
 \$\begin{aligned}
 \$U_l(\tau) = iP_l(\tau)U_l(\tau), \$P_l(\tau) \exists P_l^a(\tau)T^a\$

 \$\begin{aligned}
 \$P_l^a(\tau) = -D_l^a \begin{bmatrix} A_{\text{gauge}}(U(\tau)) \begin{bmatrix} + \eta^{\daggel}(\tau)D_l^a \begin{bmatrix} DD^{\daggel}(U(\tau)) \begin{bmatrix} + \eta^{\daggel}(\tau)D_l^a \begin{bmatrix} DD^{\daggel}(U(\tau)) \begin{bmatrix} + \eta^{\daggel}(\text{triangle}(\tau)) \begin{bmatrix} + \eta^{\daggel}(\tau)D_l^a \begin{bmatrix} DD^{\daggel}(U(\tau)) \begin{bmatrix} + \eta^{\daggel}(U(\tau)) \begin{bmatrix} + \eta^{\daggel}(\text{triangle}(\tau)) \begin{bmatrix} + \eta^{\daggel}(\text{triangle}(\text{t
- 7. Go to 2.

Conjugate Gradient algorithm $(CC^{\dagger})\eta' = \phi'$

$$Ax = b, \ A \equiv CC^{\dagger}$$
$$x_{0} = 0, \ r_{0} = b - Ax_{0}, \ p_{0} =$$
$$\alpha_{k} = \frac{(r_{k}, r_{k})}{(p_{k}, Ap_{k})} = \frac{(r_{k}, r_{k})}{(Cp_{k}, Cp_{k})}$$

$$\phi' = C_{_{PV}} \phi$$

 $\eta = C_{_{PV}}^{\dagger} \eta'$

$$r_{k+1} = r_k - \alpha_k A p_k \quad \text{If } |r_{k+1}| < \varepsilon |b|,$$

$$x_{k+1} = x_k + \alpha_k p_k \quad \text{* The most time-}$$

$$are the matrix-$$

$$c(C^{\dagger} p_k)$$

$$p_{k+1} = r_{k+1} + \beta_k p_k \quad \text{* GPU computes}$$

* The most time-consuming operations
are the matrix-vector multiplications
$$C(C^{\dagger}p_k)$$

then stop

in SP \gg DP

 r_0

CG Algorithm with Mixed Precision

- 1. $r_k = b Ax_k$
- 2. If $|r_k| < \varepsilon |b|$, then stop
- 3. Solve $At_k = r_k$ in single precision to an accuracy $\varepsilon_1 < 1$
- 4. $x_{k+1} = x_k + t_k$
- 5. Go to 1.

Proof of convergence:

Let $s_k = r_k - At_k$, $|s_k| < \varepsilon_1 |r_k|$ then $|r_{k+1}| = |b - Ax_{k+1}| = |b - Ax_k - At_k| = |s_k| < \varepsilon_1 |r_k| < |r_k|$

Hardware Model



T10/ GTX285

Global memory: # of multiprocessors: # of cores: Constant memory: Shared memory/block: # of registers/block: Warp size: Max. # threads/block: Block-size limits x,y,z: Grid-size limits in x,y: 4/1/2 Gbytes 30 240 65536 bytes 16384 bytes 16384 32 512 512, 512, 64 65535, 65535

A set of SIMT multiprocessors with on-chip shared memory.

CUDA Programming Model

A kernel is executed by a grid of thread blocks

- A thread block is a set of threads that can cooperate with each other by:
 - Sharing data through shared memory
 - Synchronizing their execution
- Threads from different blocks cannot cooperate



Features of CG Kernel for ODWF

 $(CC^{\dagger})\eta' = \phi'$ $C \equiv I - D_{w}^{oe}YX^{-1}D_{w}^{eo}YX^{-1}$

One thread takes care of all computations at each (s, x, y, z) with a loop going over all t (even/odd).

- 1-dim Grid $Nthread_X \times Nthread_Y \times Nblock = N_s N_x^3$
- 2-dim Block $Nthread_X = N_s = 16 \times (1, 2, ...)$ $Nthread_X \times Nthread_Y = 64$

Tuning the CG Kernel for ODWF

- $M \equiv YX^{-1}$ in constant memory space
- Use texture memory for link variables and vectors
- Reorder data in the device memory (coalescing; Ns threads share the same link variables)
- Reuse forward/backward data (in t) for neighboring sites
- Unroll short loops

• . .

Attaining 140/120/100 Gflops for GTX285/GTX280/T10

Salient Features of TWQCD's HMC Simulation

- HMC with Multiple Time Scale Integration and Mass Preconditioning
- Even-Odd Preconditioning for the 4D Quark Matrix
- Omelyan Integrator for the Molecular Dynamics
- Conjugate Gradient with Mixed Precision
- New Algorithm for (2+1)-flavor QCD
- Chiral Symmetry is preserved exactly with ODWF
- All Topological Sectors are sampled ergodically.
- The First Large-Scale Simulation of Unquenched Lattice QCD with Exact Chiral Symmetry using a GPU Cluster.

Technical Breakthrough

The First Large-Scale Dynamical Lattice QCD Simulation with a GPU Cluster

The lattice QCD group (TWQCD) based at NTU is the first group around the world to use a GPU cluster to perform large-scale simulations of lattice QCD with exact chiral symmetry. [T.W. Chiu et al., PoS (LAT2009) 034, arXiv:0911.0529].

Currently, there are only 3 groups (RBC/UKQCD, JLQCD, and TWQCD) around the world who are capable to perform large scale simulation of lattice QCD with exact chiral symmetry.

Research Highlights

> Zero Temperature QCD $(N_f = 2)$



T.W. Chiu, GPU Computing Seminar, January 22, 2010

Research Highlights (cont)



Fitting data to $f_{\pi} = f(1 - x \ln x) + c_4 x \implies f a = 0.091(6)$ Using f = 131 MeV as input $\implies a^{-1} = 1.45(9)$ GeV

Research Highlights (cont)

Topological structure of the QCD vacuum

- The vacuum (ground state) of QCD constitutes various quantum fluctuations, which are the origin of many interesting and important nonperturbative physics
- In QCD, each gauge configuration possesses a well-defined topological charge Q with integer value.
- It is important to determine the topological charge fluctuations in the QCD vacuum. One of the outstanding problems is to identify the vacuum fluctuations which are crucial to the mechanism of color confinement

Quantum fluctuations in the QCD vacuum



Research Highlights (cont)

Finite Temperature QCD

The nature of finite temperature phase transition in QCD is relevant to the evolution of the early universe from the **quark-gluon plasma** to the **hadrons**.

However this issue has remained controversial for many years. The Wuppertal group finds the deconfining transition [170(7) MeV] and the chiral transition [146(5) MeV] are at widely separated temperatures. On the other hand, the Brookhaven/Bielefeld collaboration claims that both temperatures coincide at 196(3) MeV. Since both groups used the staggered fermion, it is vital to resolve this issue in the framework of lattice QCD with exact chiral symmetry. We are in a good position to resolve this problem.



Big Bang Timeline

<u>Time</u>	<u>Era</u>	<u>Temperature</u>	Characteristics of the Universe
		_	infinitely small, infinitely dense
0 to 10 ⁻⁴³ s	Big Bang	infinite	Primeval fireball
			1 force in nature - Supergravity
			Earliest known time that can be
10 ⁻⁴³ s	Planck Time	10 ³² K	described by modern physics
			2 forces in nature, gravity, GUT
			3 forces in nature, gravity, strong
10- ³⁵ c	End of CUT	1027 K	nuclear, electroweak
10 5		10 K	Quarks and leptons form
			(along with their anti-particles)
10 ⁻³⁵ to 10 ⁻³³ s	Inflation	10 ²⁷ K	Size of the Universe drastically
10 10 10 3	mation	10 K	increased, by factor of 10^{30} to 10^{40}
			4 fundamental forces in nature,
10 ⁻¹² s	End of unified forces	10¹⁵ K	protons and neutrons start
			forming from quarks
10-7 a	Hoorn Dontialo	1014 V	proton, neutron production
10 ' S	neavy Particle	10 K	in full swing
10 ⁻⁴ s	Light particle	10 ¹² K	electrons and positrons form
100 s (a few	Nucleosynthesis or	109 107 V	helium, deuterium, and a few other
minutes)	Nucleosynthesis era	$10^{2} - 10^{2} \text{ K}$	elements form
			Matter and radiation seperate
380,000 years	Recombination (Decoupling)	3000 K	End of radiation domination, start of
			matter domination of the Universe
500 million urs	Calaxy formation	10 V	galaxies and other large structures
500 minion yrs	Galaxy IOI IIIation	10 K	form in the universe
14 billion years	Now	212	You are reading this table, that's what's
or so	INO VV	JI	happening.

Conclusion and Outlook

- GPU has emerged as a revolutionary device for lattice QCD as well as any computational science.
 It is crucial to TWQCD's dynamical DWF project.
- Currently, there are many lattice QCD groups around the world building large GPU clusters dedicated to QCD. This will have significant impacts to lattice QCD, leading to new discoveries in the strong interaction physics.